EFFECT OF ELECTRICAL RESISTIVITY ON THE DAMPING OF SLOW SAUSAGE MODES

Motivation

Dependence on magnetic Reynolds number of damping-time-to-period ratio of sausage modes in a solar photospheric pore studied numerically by [2] in the resistive $(\eta \neq 0)$ MHD framework:



- Region of $R_m > 10^7$: damping almost independent of $\eta \Rightarrow$ mainly due to resonant absorption in cusp continuum
- Region of $R_m < 3 \cdot 10^4$: damping linearly dependent on $\eta \Rightarrow$ mainly due to resistive effects
- Green region of $3 \cdot 10^4 < R_m < 10^7$: intermediate regime where both electrical resistivity and resonant absorption are important for damping

Goal: to explain/confirm this behavior through an analytical model



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Straight cylinder aligned with equilibrium magnetic field, circular basis and discontinuous boundary (l = 0), as a model for the photospheric pore (figure modified from [1]). Inside (index "i") and outside (index "e") of cylinder are two distinct uniform plasmas.

Analytical derivations

• Every perturbed quantity is expressible in terms of $\nabla \cdot \xi = R(r) \exp \{i (k_z z - \omega t)\}$. The following solution is found for R:

$$R(r) = \begin{cases} C_1 \ J_0(\kappa_{-,i}r) + C_3 \ J_0(\kappa_{+,i}r) & \text{if } r < a \\ C_2 \ H_0(\kappa_{-,e}r) + C_4 \ H_0(\kappa_{+,e}r) & \text{if } r > a \end{cases},$$

where J_0 is Bessel function of first kind and order 0, and H_0 means either $H_0^{(1)}$ or $H_0^{(2)}$ (Hankel functions of order 0), depending on the sign of the imaginary part of its argument.

• κ_{-} represents the wave part of the solution, which is complex because of damping from resistivity. It can be thought of as the radial wavenumber

 κ_+ represents Hartmann layer (electromagnetic boundary layer with strong gradients which gets thinner as resistivity η becomes smaller)

- With proper boundary conditions, complex dispersion relation for ω can be derived
- Long wavelength limit formula of dispersion relation also available

Results

• Solving dispersion relation numerically allows to compare **damping-time-to period ratio** (τ/P) as a function of magnetic Reynolds number (R_m) from analytical model ("ana.") mod.") of this research with numerical calculations ("num. cal.") from [2] code adapted to infinitely small transition layer $(l \approx 0)$:







• In the long wavelength limit $(k_z a \rightarrow 0)$, analytical results produce simpler formula for damping rate. Since ω close to internal cusp frequency ($\omega_{\rm Ci}$) in this limit, writing

we find:

$$\omega^2 = \omega_{\rm Ci}^2 (1 + \delta)$$

$$\delta \propto \sqrt{r}$$

• Following figure is graph of **damping time-to-period ratio** in function of $k_z a$ for long wavelength limit (in logarithmic scale):



In the long wavelength limit, the resistive damping thus vanishes.

References & acknowledgements

- [1] Arregui, I. et al. "Resonantly damped fast MHD kink modes in longitudinally stratified tubes with thick non-uniform transitional layers". In: Astronomy & Astrophysics 441.1 (2005), pp. 361-370
- [2] Chen, S. et al. "Damping of Slow Surface Sausage Modes in Photospheric Waveguides". In: *The* Astrophysical Journal 868.1 (2018), p. 5.

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