

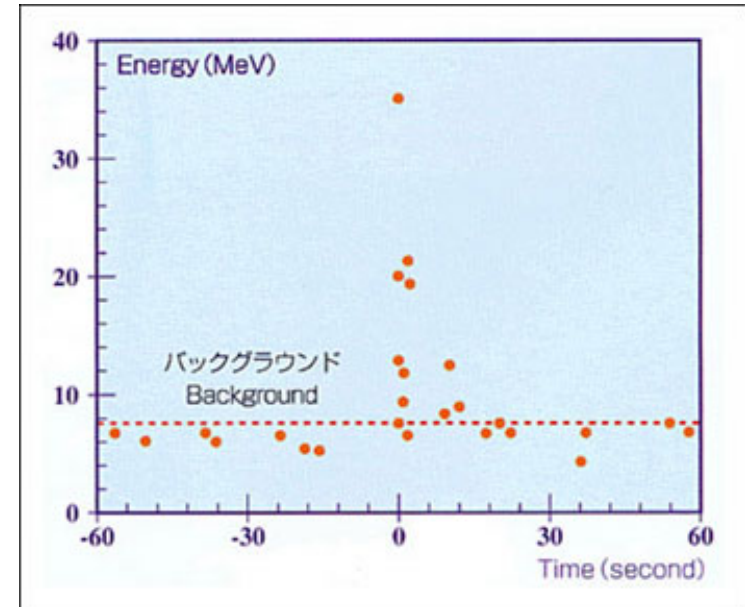
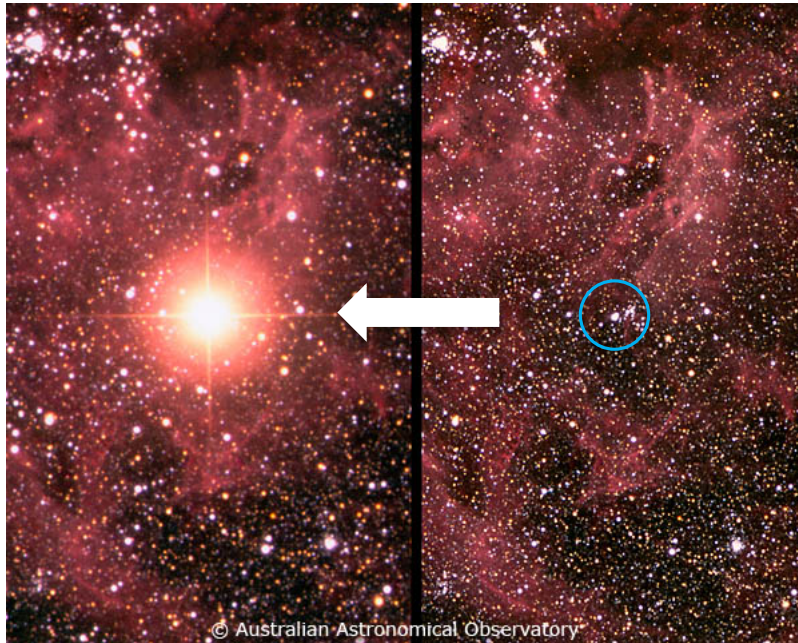
# Collective neutrino oscillations in core-collapse supernovae

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The University of Tokyo

Workshop on core-collapse supernova explosions and related physics  
6<sup>th</sup> August 2019 at University of Exeter in UK

- Introduction
  - Collective Neutrino Oscillation (CNO)
  - BULB model
  - Beyond BULB model
    - Recent trials for symmetry breakings
- My recent work
  - Neutrino halo effects in iron-CCSNe

# Introduction



Kamiokande-II and IBM detector observed 20 neutrino events from SN1987A in LMC.

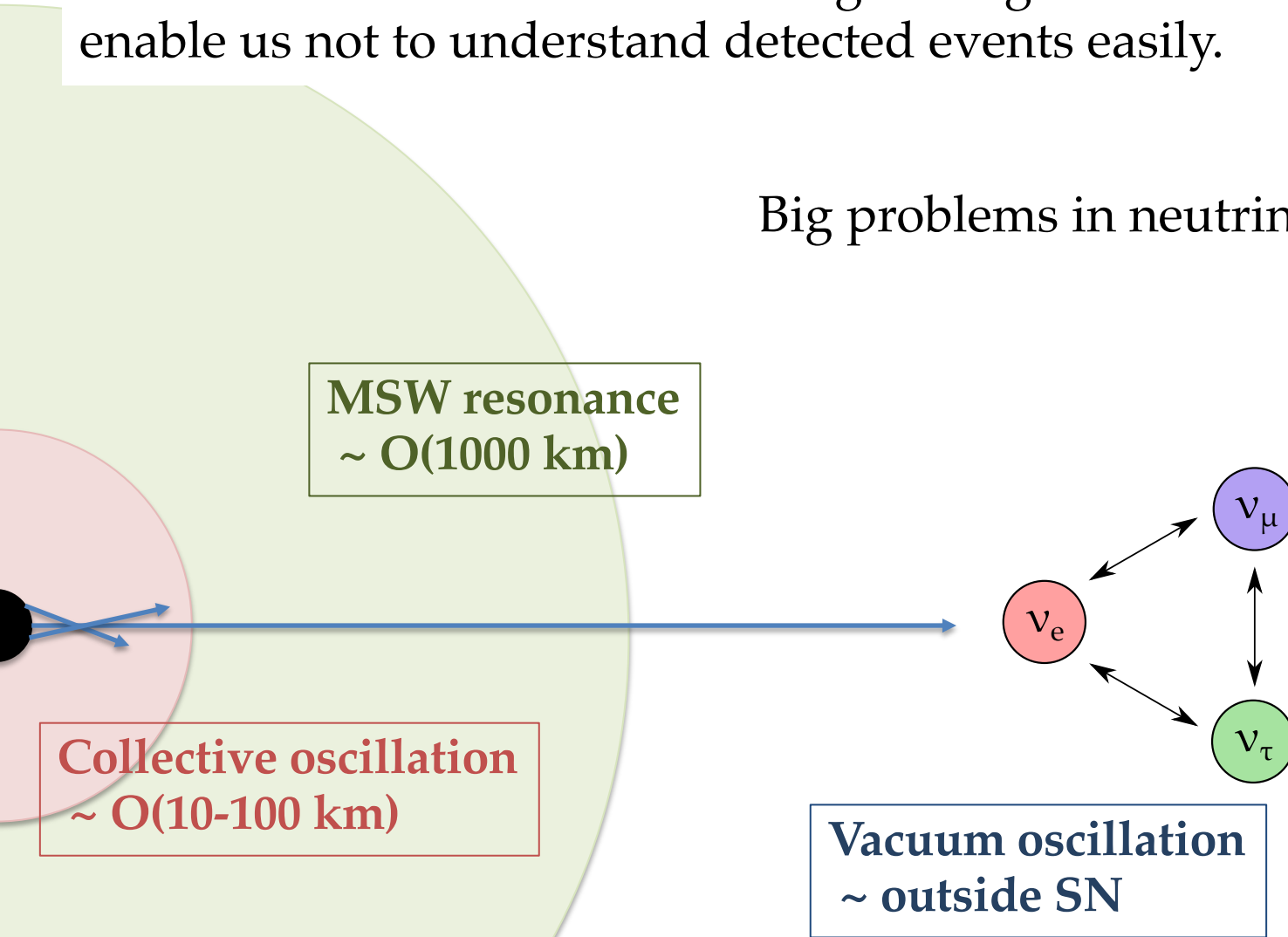
This event taught us much neutrino information.  
average energy, luminosity, emission duration.

What will next galactic supernova bring...?

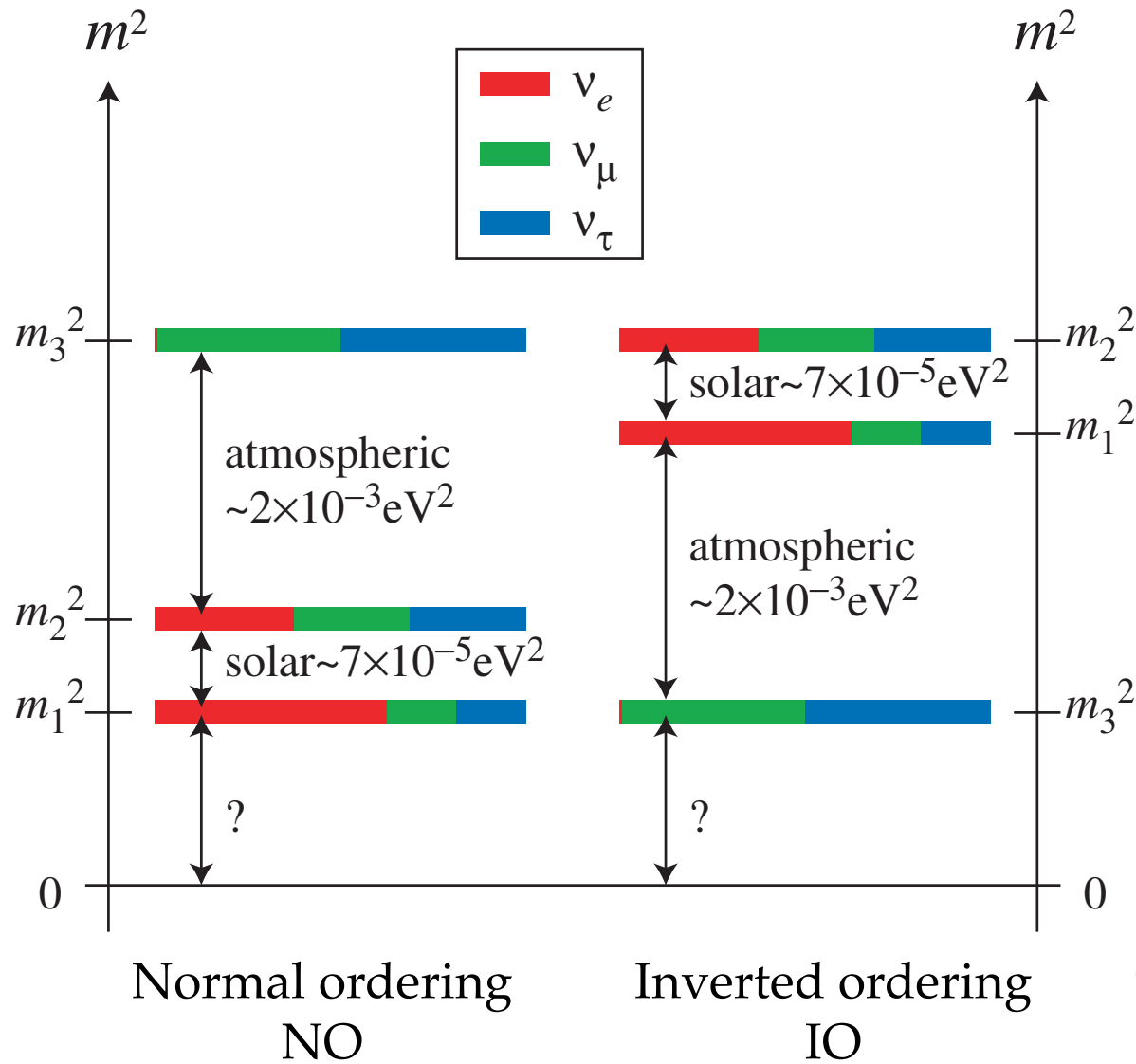
# Neutrino oscillations in CCSNe

Neutrino oscillation causes mixing among neutrino flavors, and enable us not to understand detected events easily.

Big problems in neutrino physics.



# Neutrino mass ordering



# Neutrino oscillation

Neutrino oscillations have three types in supernovae.

$$i (\partial_t + \mathbf{v} \cdot \nabla) \rho_\nu = [H_{\text{tot}}, \rho_\nu]$$

$$\rho_\nu = \begin{pmatrix} f_{\nu_e} & f_{\langle \nu_e | \nu_x \rangle} \\ f_{\langle \nu_x | \nu_e \rangle} & f_{\nu_x} \end{pmatrix}$$

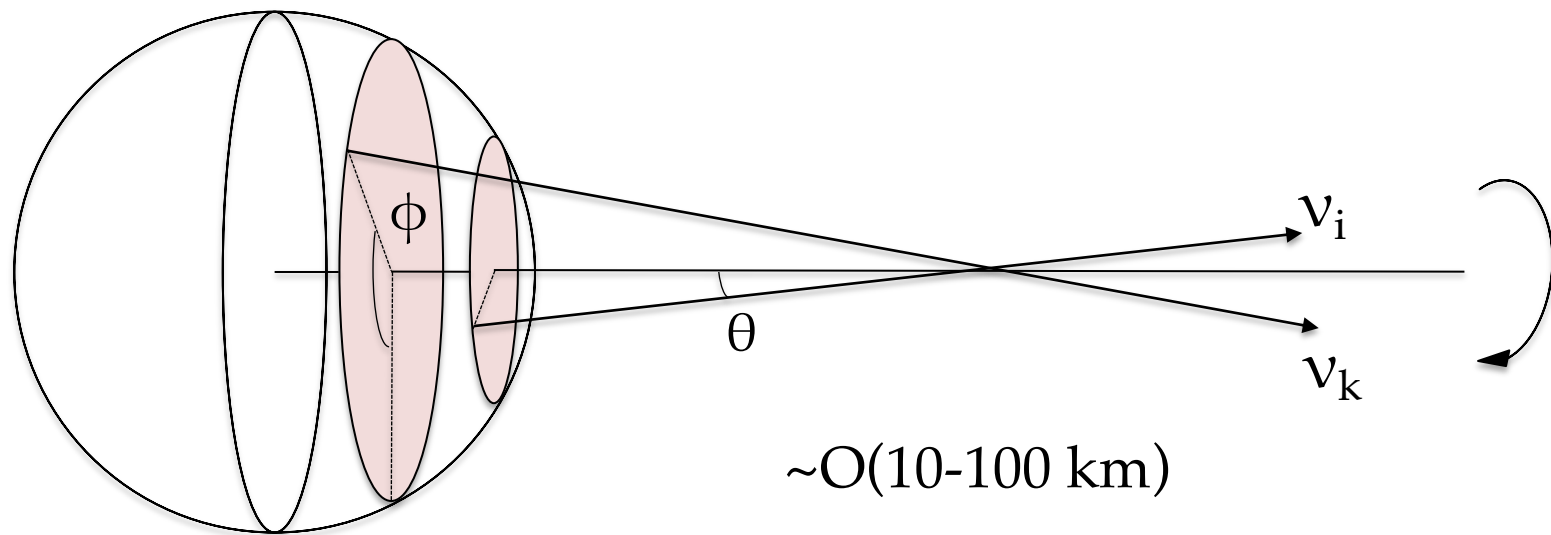
$$H_{\text{tot}} = \underbrace{\frac{M^2}{2E}}_{\text{Vacuum oscillation (Neutrino mass)}} + \underbrace{\sqrt{2}G_F n_e}_{\text{Matter oscillation (Electron density)}} + \underbrace{\sqrt{2}G_F \int d\Gamma' (1 - \mathbf{v} \cdot \mathbf{v}') (\rho_\nu - \bar{\rho}_\nu)}_{\text{Collective neutrino oscillation (CNO) (v-v forward scattering)}}$$

Vacuum oscillation  
(Neutrino mass)

Matter oscillation  
(Electron density)

Collective neutrino oscillation (CNO)  
(v-v forward scattering)

# Collective neutrino oscillation



PNS

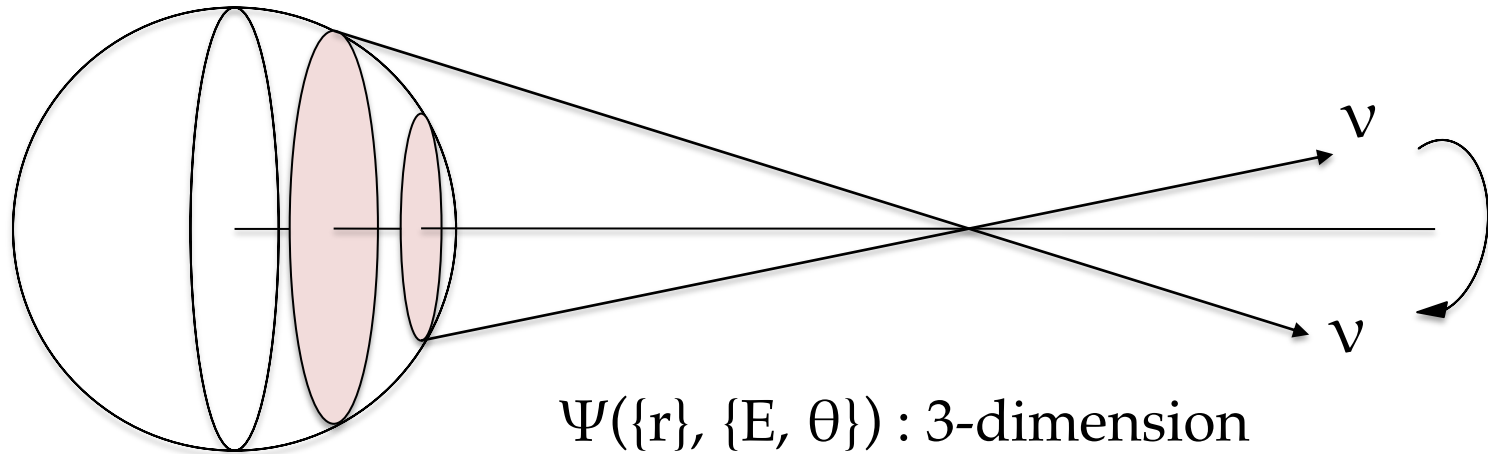
$\Psi(t, \{r, \Theta, \Phi\}, \{E, \theta, \phi\}) : 7\text{-dimension}$

$\sim 10^{58}$  neutrinos are emitted from a SN in  $\sim 10$  seconds.

Dense neutrino gas causes neutrino-neutrino interaction near the proto-neutrino star.

Neutrino self-interaction induces collective flavor conversions.

# Collective neutrino oscillation



$\Psi(\{\mathbf{r}\}, \{E, \theta\}) : 3\text{-dimension}$

$N_{\text{flavor}} \times N_E \times N_\theta = O(10^7)$  ODE systems

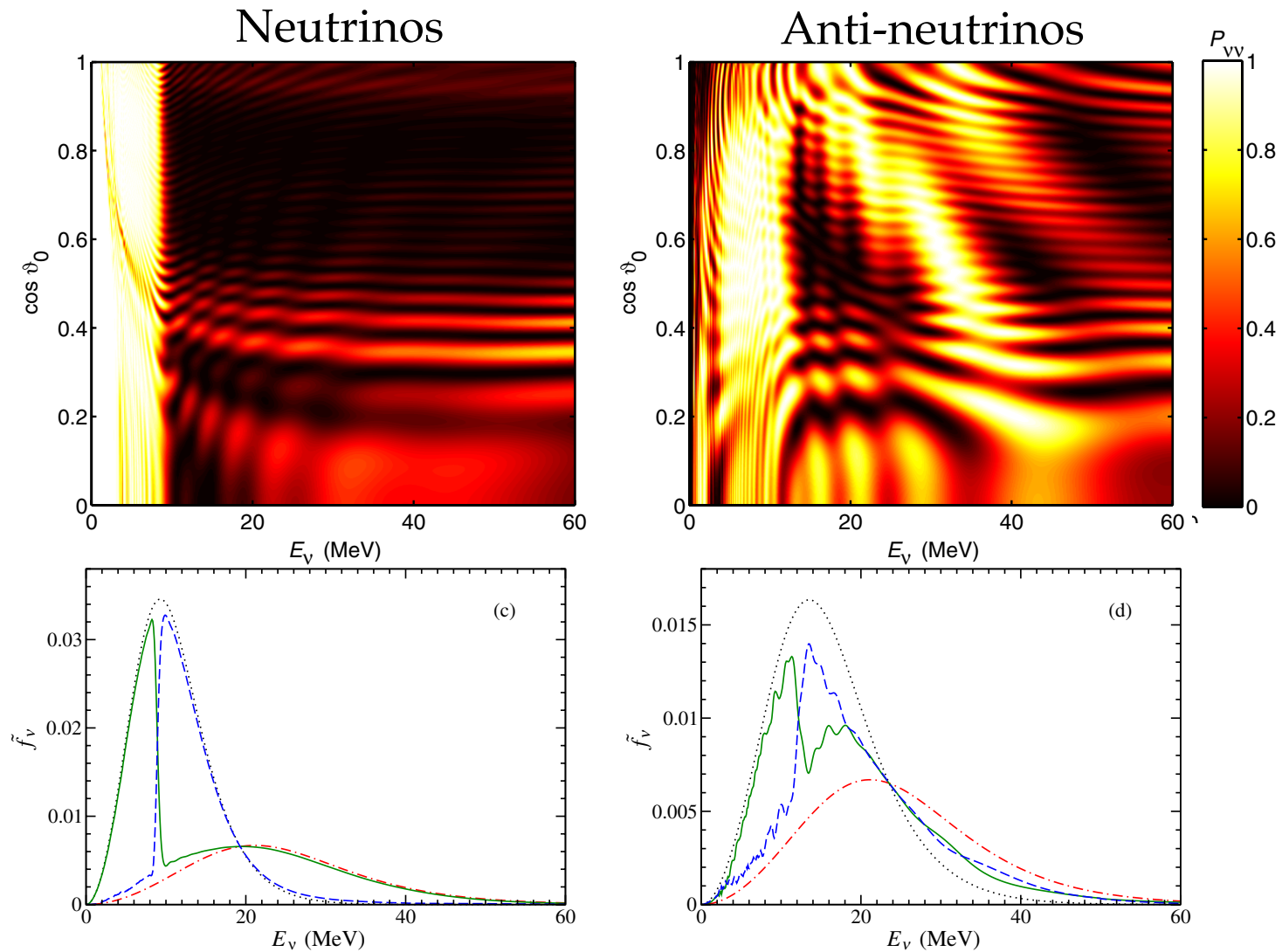
The “**BULB model**” to simplify the neutrino self-interaction.

- Homogeneous, Isotropic, Stationary emission
- Axial-symmetry in direction
- 1-dimensional SN model
- Neutrinos propagate free-streamingly.

(Duan+ PRD 2006, Dasgupta+ PRD 2008)



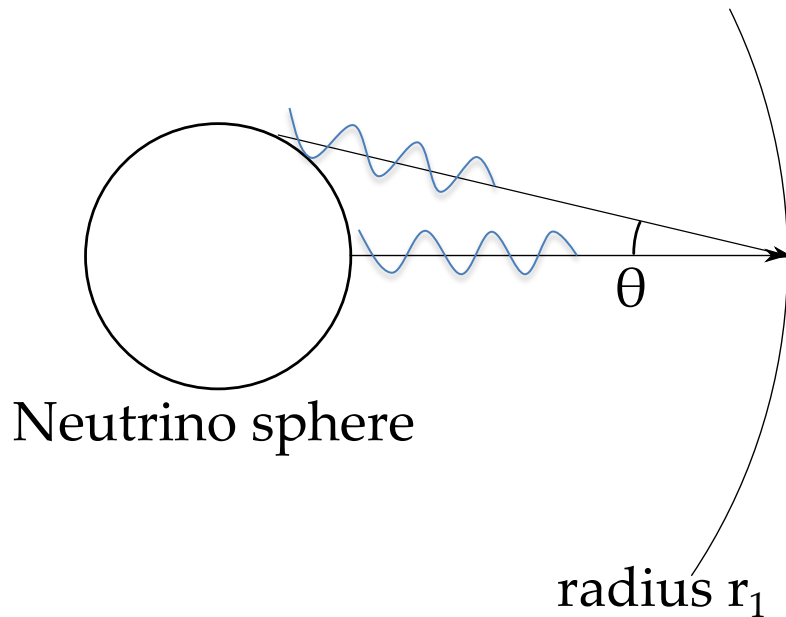
# Spectral splits



(Duan+ PRL, PRD 2006)

Flavor conversions above a critical energy

# Collective effects vs. Matter effects



(Esteban-Pretel+ PRD 2008)

Neutrino self-interaction maintains coherence among different neutrino trajectories.

On the other hand,

Different neutrino trajectories give the different effective potential from background electrons.  
This causes phase dispersion.

- Matter effects  $>$  neutrino self-interaction  
→ Collective neutrino oscillation can be suppressed.
- Matter effects  $<$  neutrino self-interaction  
→ Collective flavor transformation occurs.

# Equation of motion

$$i\partial_t \rho_i = \left[ \underbrace{+U \frac{M^2}{2E_\nu} U^\dagger}_{\text{Vacuum oscillation}} + \underbrace{\sqrt{2}G_F N_e}_{\text{Matter oscillation}} + \underbrace{\sqrt{2}G_F \int d\Gamma' (1 - \cos \theta_{ik}) (\rho'_{\nu_k} - \bar{\rho}'_{\nu_k})}_{\text{Collective neutrino oscillation}}, \rho_i \right]$$

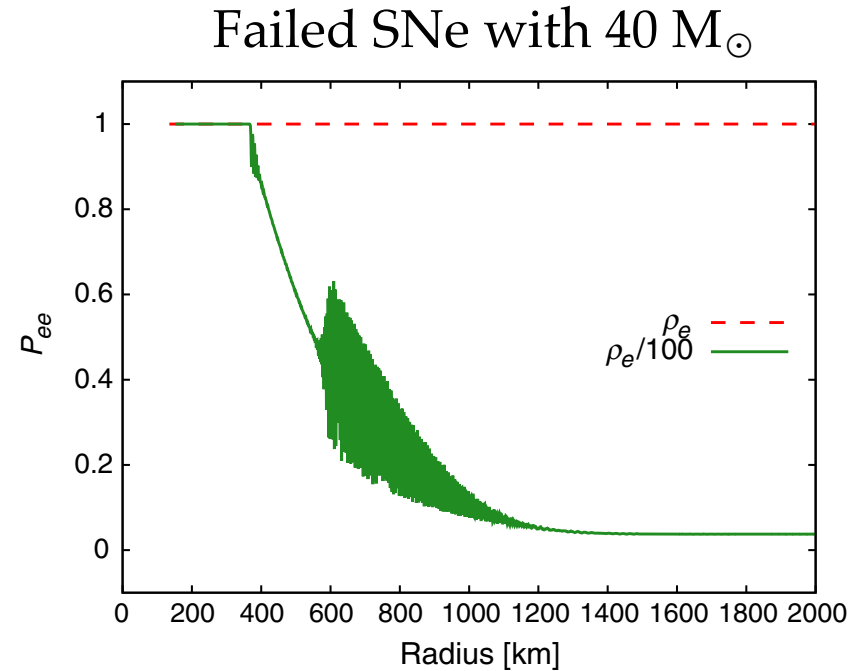
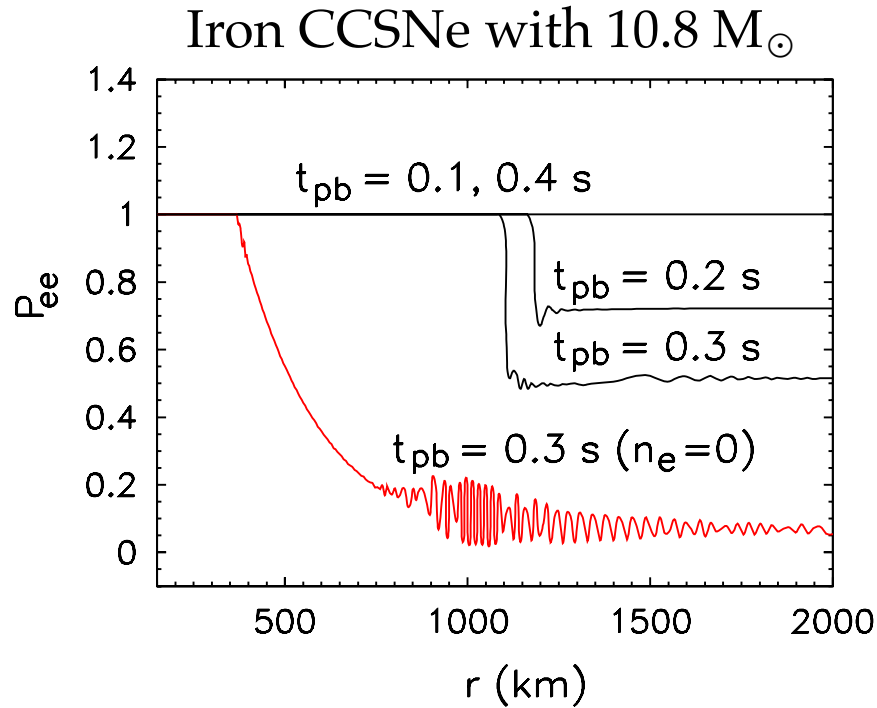
Vacuum oscillation  
Mass difference

Matter oscillation  
Density profile

Phase dispersion  
vs.  
Phase synchronization

Collective neutrino oscillation  
Neutrino trajectory & Neutrino spectra

# Matter suppression

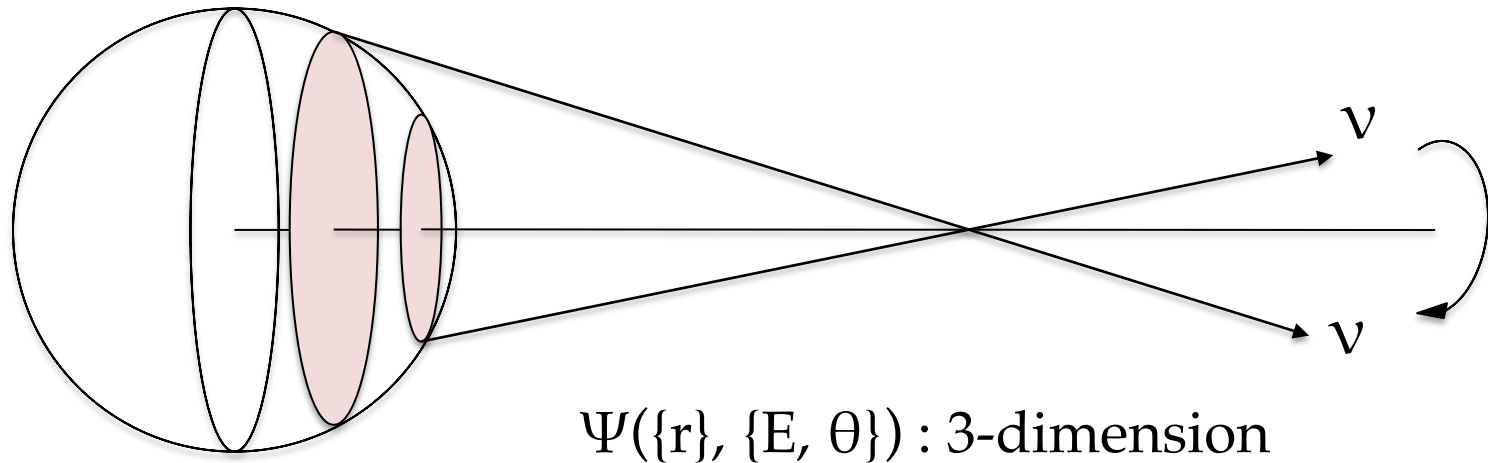


(Chakraborty+ PRD 2011, Zaizen+ PRD 2018)

CNO is suppressed in SNe with low mass progenitor at early time, especially at accretion phase.

CNO is completely suppressed in SNe with heavy progenitor at all time.

# Limitation of BULB model?



**The BULB model requires many symmetries & assumptions.**

- Stationary emission
- 1-dimensional SN model
- Directional axial-symmetry
- Homogeneous, Isotropic emission
- Neutrinos propagate free-streamingly.

# Symmetry breaking

3-dimensional EoM --> 7-dimensional EoM

$$iv_r \partial_r \rho_\nu = [H_{\text{tot}}, \rho_\nu]$$

$$i(\partial_t + v_r \partial_r + \mathbf{v}_T \cdot \nabla_T) \rho_\nu = [H_{\text{tot}}, \rho_\nu]$$

$$\rho_\nu (t, \{r, \Theta, \Phi\}, \{E, \theta, \varphi\}) : 7 \text{ dimension}$$

## The BULB model – three types of symmetries

- Stationary emission → Temporal instability
- 1-dimensional SN model → Homogeneous instability
- Axial-symmetry in direction → Multi-azimuthal angle instability

# Stability analysis

$$i(\partial_t + v_r \partial_r + \mathbf{v}_T \cdot \nabla_T) \rho_\nu = [H_{\text{tot}}, \rho_\nu]$$

$$\rho_\nu = \begin{pmatrix} f_{\nu_e} & f_{\langle \nu_e | \nu_x \rangle} \\ f_{\langle \nu_x | \nu_e \rangle} & f_{\nu_x} \end{pmatrix} \quad f_{\langle \nu_e | \nu_x \rangle} \quad \begin{array}{l} \text{Initially this is zero.} \\ \text{No flavor evolution} \end{array}$$

Off-diagonal components of density matrix express the transitions from one flavor to another.

We can investigate unstable modes to evaluate the growth rate of transition term.

$$f_{\langle \nu_e | \nu_x \rangle} \propto e^{-i\Omega r}$$

If  $\text{Im } \Omega < 0$ , we think this is unstable at a radius  $r$ .

# Stability analysis

$$i(\partial_t + v_r \partial_r + \mathbf{v}_T \cdot \nabla_T) \rho_\nu = [H_{\text{tot}}, \rho_\nu]$$



Linearize

$$\left[ \frac{-\omega + \bar{\lambda} - p - \mathbf{v}_T \cdot \mathbf{k}}{v_r} - \Omega_{p,\mathbf{k}} \right] Q_{p,\mathbf{k}} = \frac{\mu}{v_r} \int d\Gamma' (1 - \mathbf{v} \cdot \mathbf{v}') g' Q'_{p,\mathbf{k}}$$

$$f_{\langle \nu_e | \nu_x \rangle} = \int dp d\mathbf{k} e^{-i(pt + \mathbf{v}_t \cdot \mathbf{k})} Q_{p,\mathbf{k}} e^{-i\Omega_{p,\mathbf{k}} r}$$

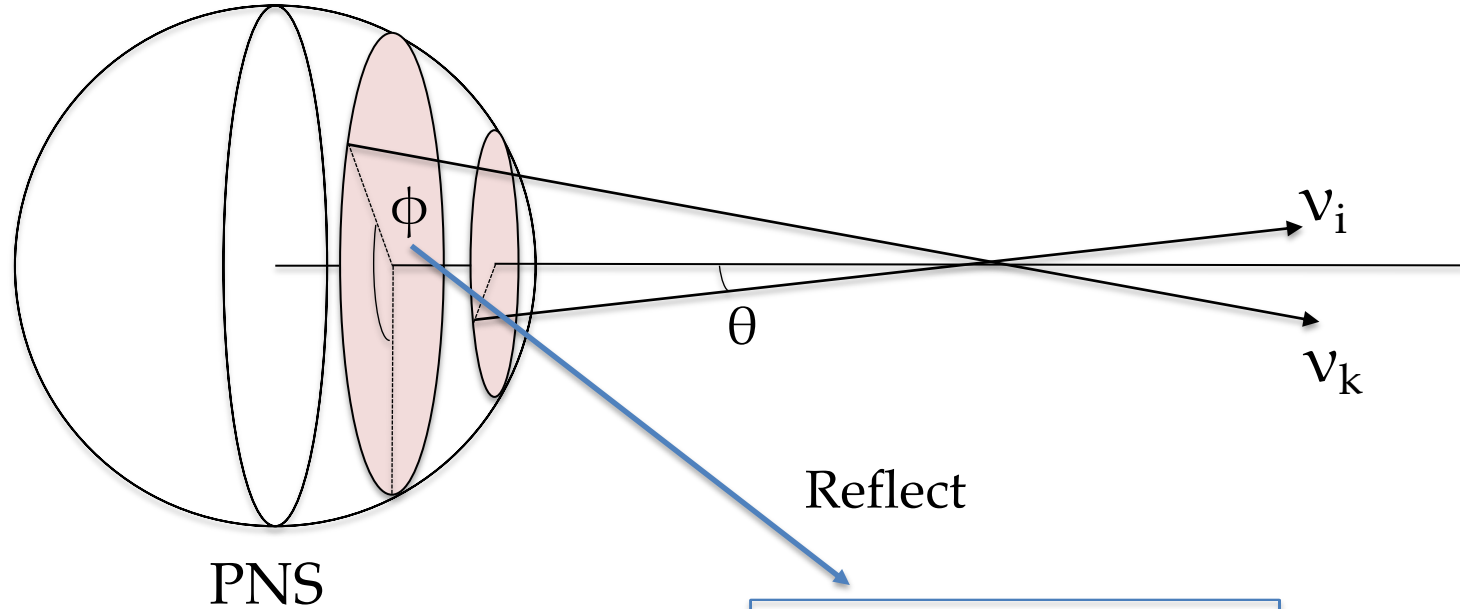
We solve this linearized equation.

$p$  &  $\mathbf{v} \cdot \mathbf{k}$  arises from derivative term  $\partial_t$  &  $\mathbf{v} \cdot \nabla$ .

BULB model ignores  $p$  &  $\mathbf{v} \cdot \mathbf{k}$  terms.



# Axial-symmetry breaking

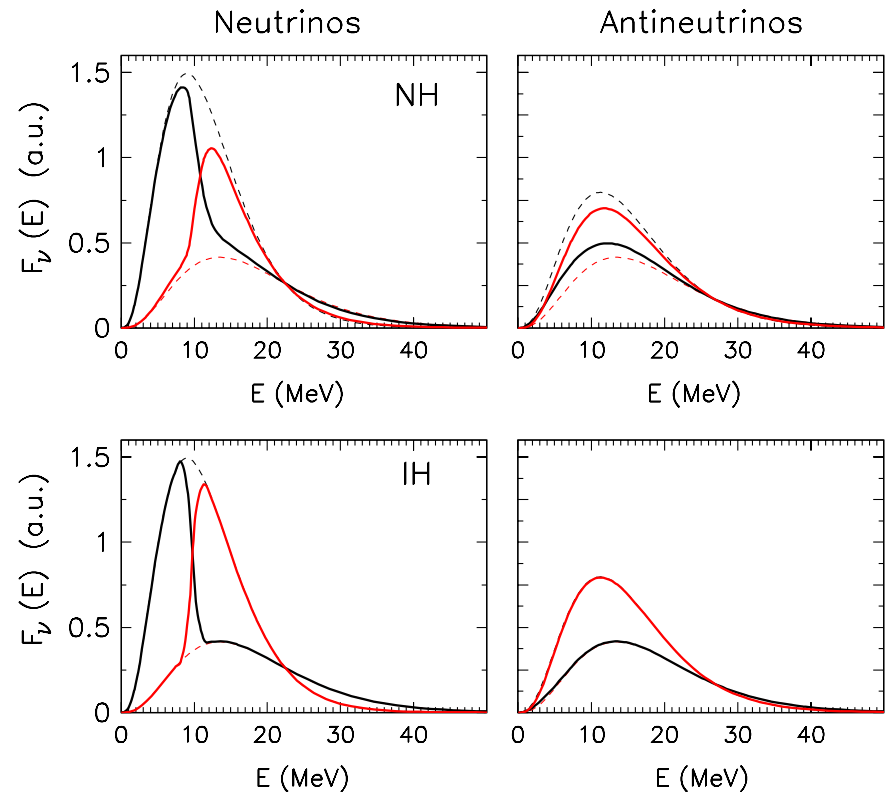
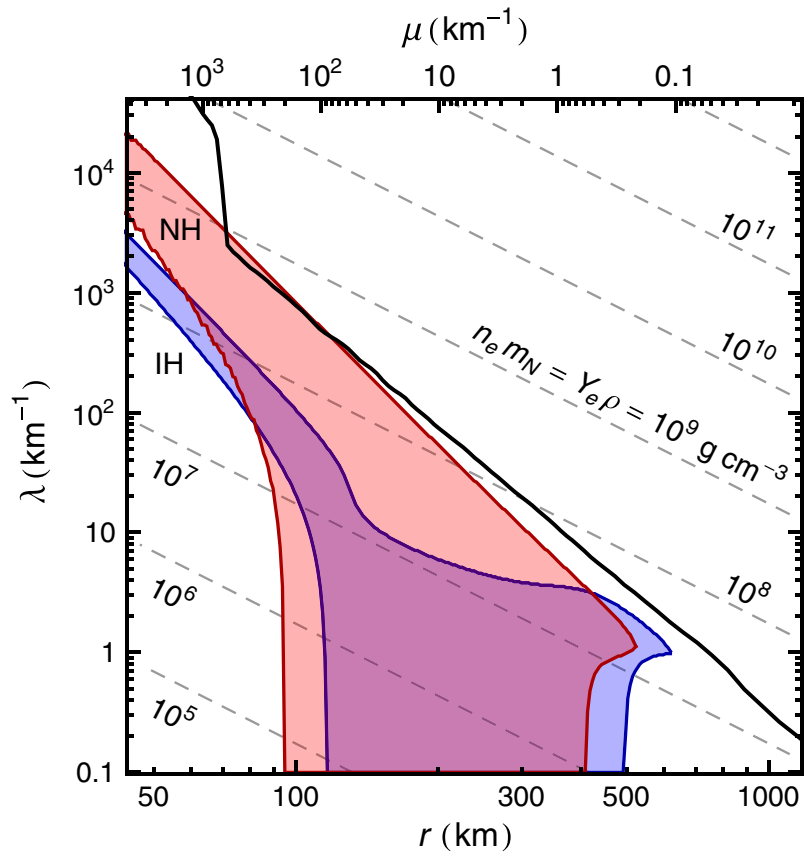


$$H_{\nu\nu} = \int d\Gamma' \left( \frac{1 - v_{r,u}v_{r,u'} - \sqrt{uu'} \frac{R_\nu^2}{r^2} \cos(\varphi - \varphi')}{v_{r,u}v_{r,u'}} \right) (\rho_\nu - \bar{\rho}_\nu)$$

Add an azimuthal-angle term in neutrino trajectories.

This requires angular binning of  $N_\phi \leq 100$ , less than  $N_\theta \sim 1000$  ?

# Axial-symmetry breaking

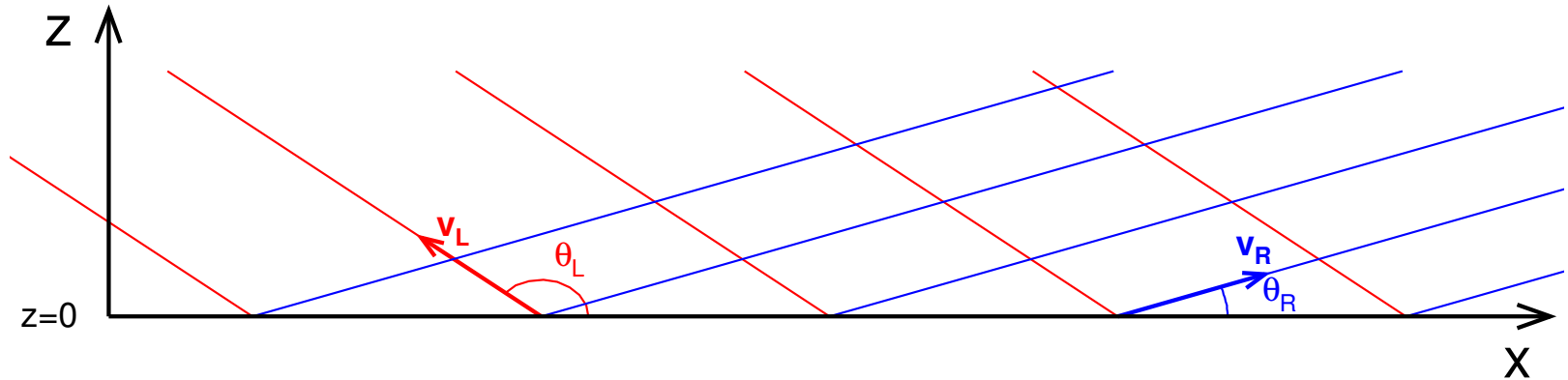


(Raffelt+ PRL 2013, Chakraborty+ PRD 2014)

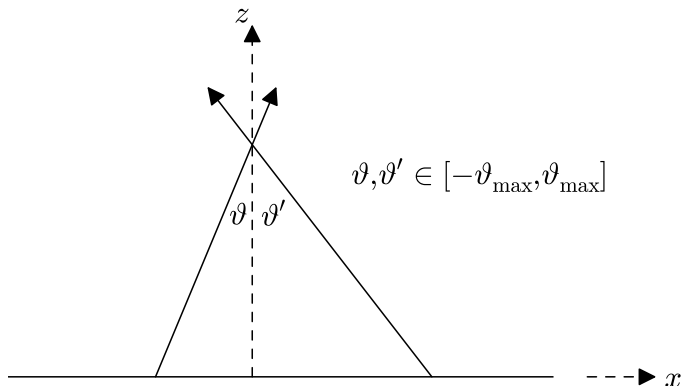
This axial-symmetry breaking indicates that CNO is more unstable in the normal mass ordering than in the inverted ordering.

This feature is distinct from the BULB model!

# Spatial-symmetry breaking



Neutrinos are emitted from the  
``**neutrino line**'', not the neutrino sphere.

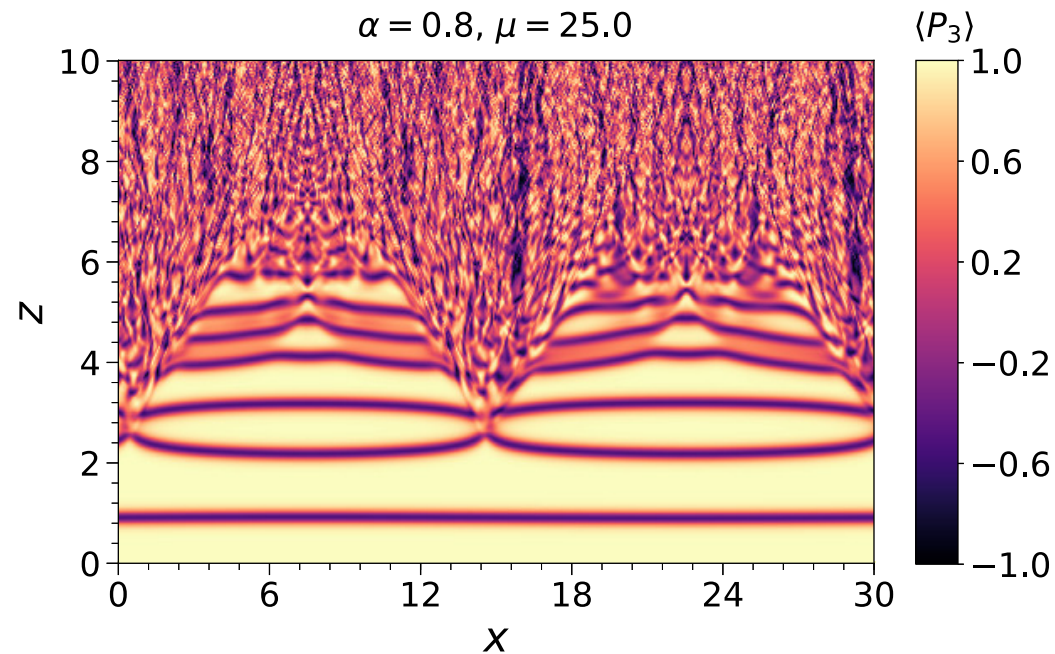


Simply studies often use a two-beam  
model, which has a single angle.

(Of course, multi-angle beam model is  
more realistic.)

# Spatial-symmetry breaking

$$i (v_r \partial_r + \mathbf{v}_T \cdot \nabla_T) f_{\langle \nu_e | \nu_x \rangle}$$



(Martin+ arXiv 2019)

Two-dimensional, but two-beam line model ( $\theta = \pi/4$ ).

Large-scale structures along  $x$ -direction begin to break down into small-scale structures as  $z$  increases.

Small-scale structures require much fine  $x$ -binning ( $N_x = 60,000$ )...

# Temporal instability

$$i(\partial_t + v_r \partial_r + \mathbf{v}_T \cdot \nabla_T) \rho_\nu = [H_{\text{tot}}, \rho_\nu]$$

Linearize  $f_{\langle \nu_e | \nu_x \rangle} \propto e^{-ipt - i\Omega r}$

$$\left[ \frac{-\omega + \bar{\lambda} - p}{v_r} - \mathbf{v}_T \cdot \mathbf{k} - \Omega_{p,\mathbf{k}} \right] Q_{p,\mathbf{k}} = \frac{\mu}{v_r} \int d\Gamma' (1 - \mathbf{v} \cdot \mathbf{v}') g' Q'_{p,\mathbf{k}}$$

$p$  arises from time derivative  $\partial t$ . This expresses a pulsating mode.

Pulsation mode  $p$  can cancel density term  $\lambda$ .

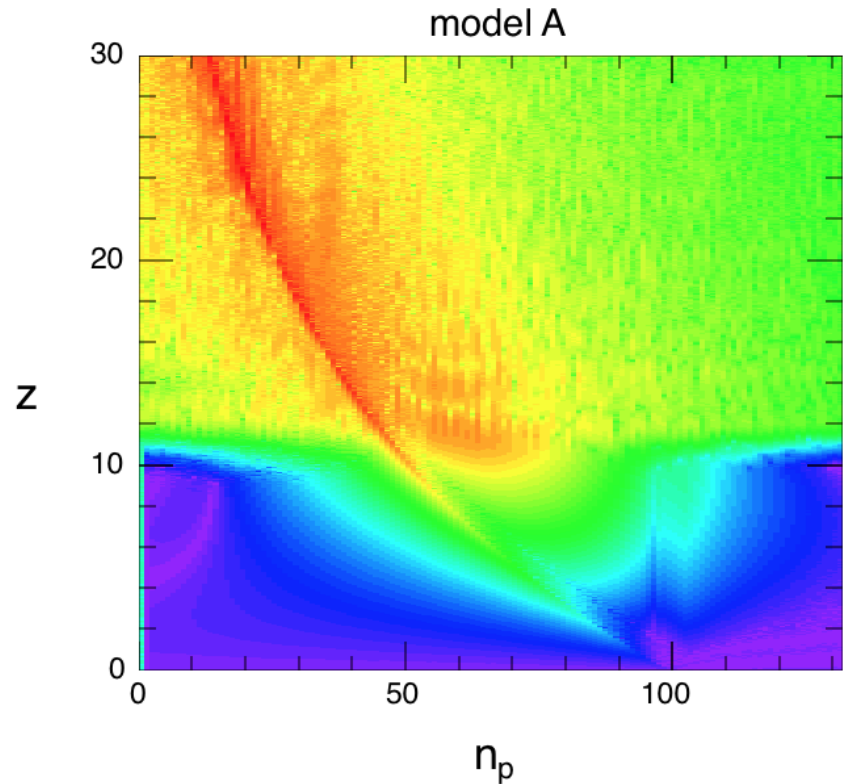
If  $\lambda$  is large enough to suppress CNO completely, instability with pulsation mode  $p \sim \lambda$  can grow up exponentially.

# Temporal instability

$$p = \frac{n_p}{100} \bar{\lambda}$$

$$\bar{\lambda} - p = 0 \rightarrow n_p \sim 100$$

$$\lambda = \lambda_0 \times \exp(-z/\tau_\lambda)$$



(Capozzi+ JCAP 2016)

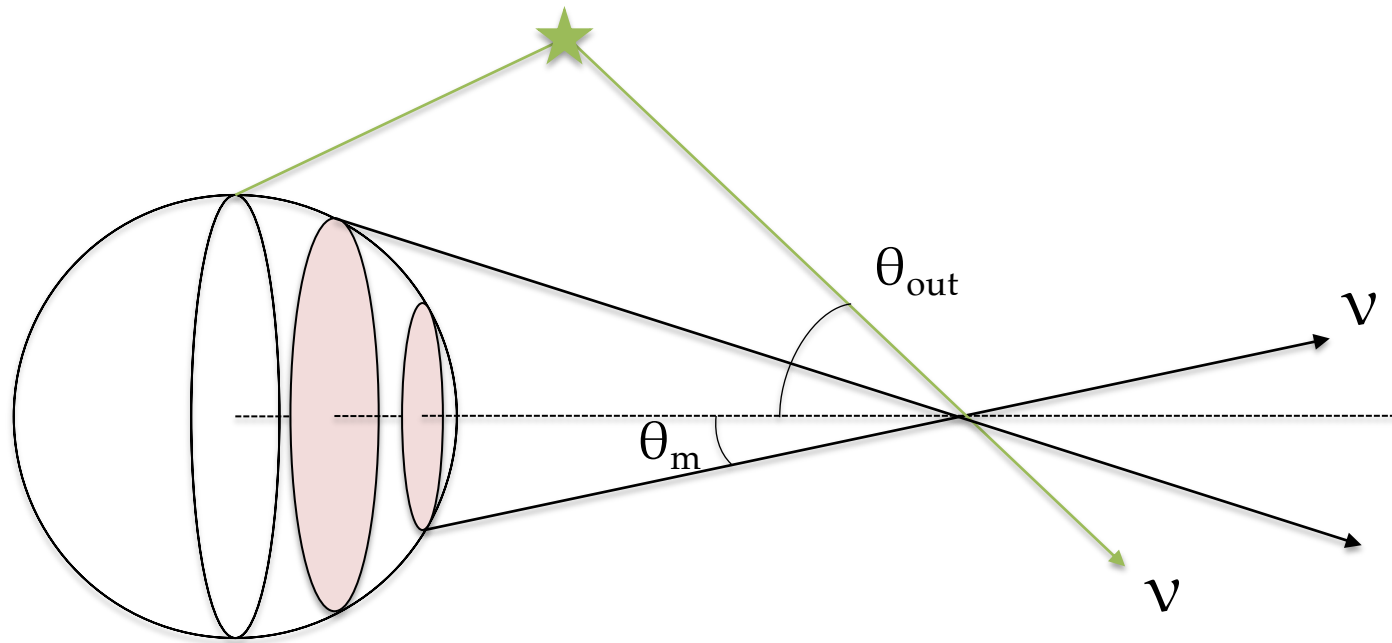
Initially modes at  $n_p \sim 100$  are excited.

Thereafter excited modes are shifted into slower pulsating modes as density declines ( $z$  decreases).

Possibility that flavor conversion can occur inside the shock wave at accretion phase?

# Neutrino halo effect on collective neutrino oscillation in iron core-collapse supernova

# Neutrino halo effect



$$(1 - \cos \theta \cos \theta')$$

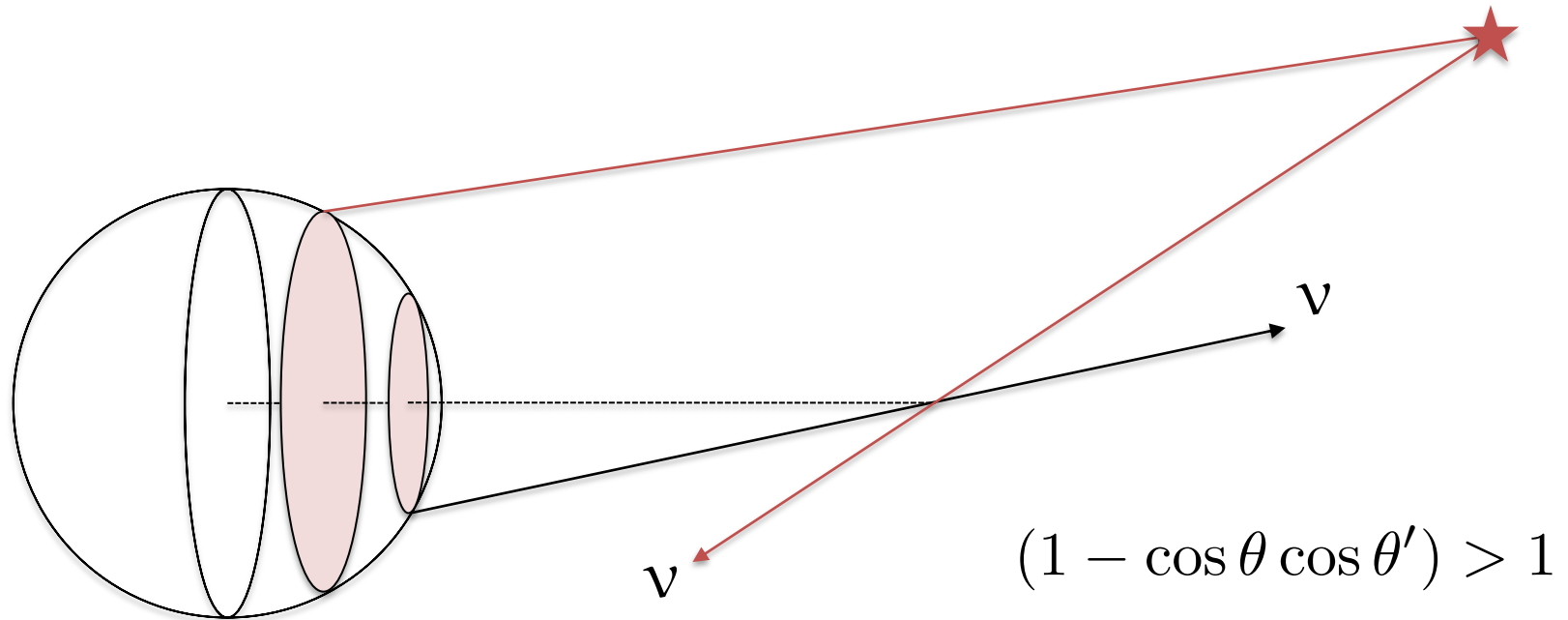
Some fraction of neutrinos emitted from neutrino sphere can experience a direction-changing scattering.

These neutrinos have wider-intersection angles, and can have impact on CNO.

(Cherry+ PRL 2012)



# Inward neutrino flux



Inward-scattered neutrinos destroy the bulb model.

We have to simulate CNO with iteration. But this is still a static case.

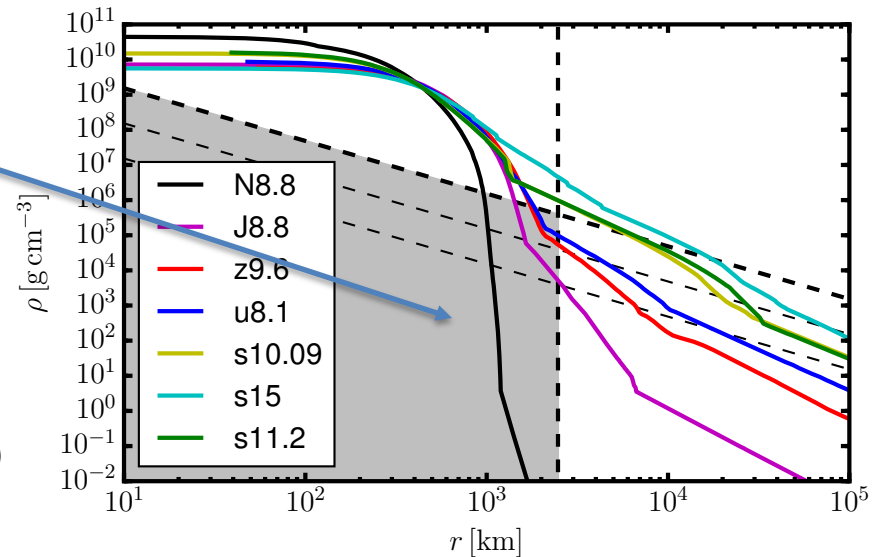
In order to calculate CNO within the bulb framework, inward flux needs to be ignored, compared with outward flux.

# Inward neutrino flux

Previous study (Cherry+ PRD 2013) employed an O-Ne-Mg SN model.

- It has a steep density gradient.

(Muller 2016)



Neutrino halo is produced by neutrino-nucleon / nucleus scattering.

- After density gradient, scattering becomes much smaller.
- Inward scattering is also much less.

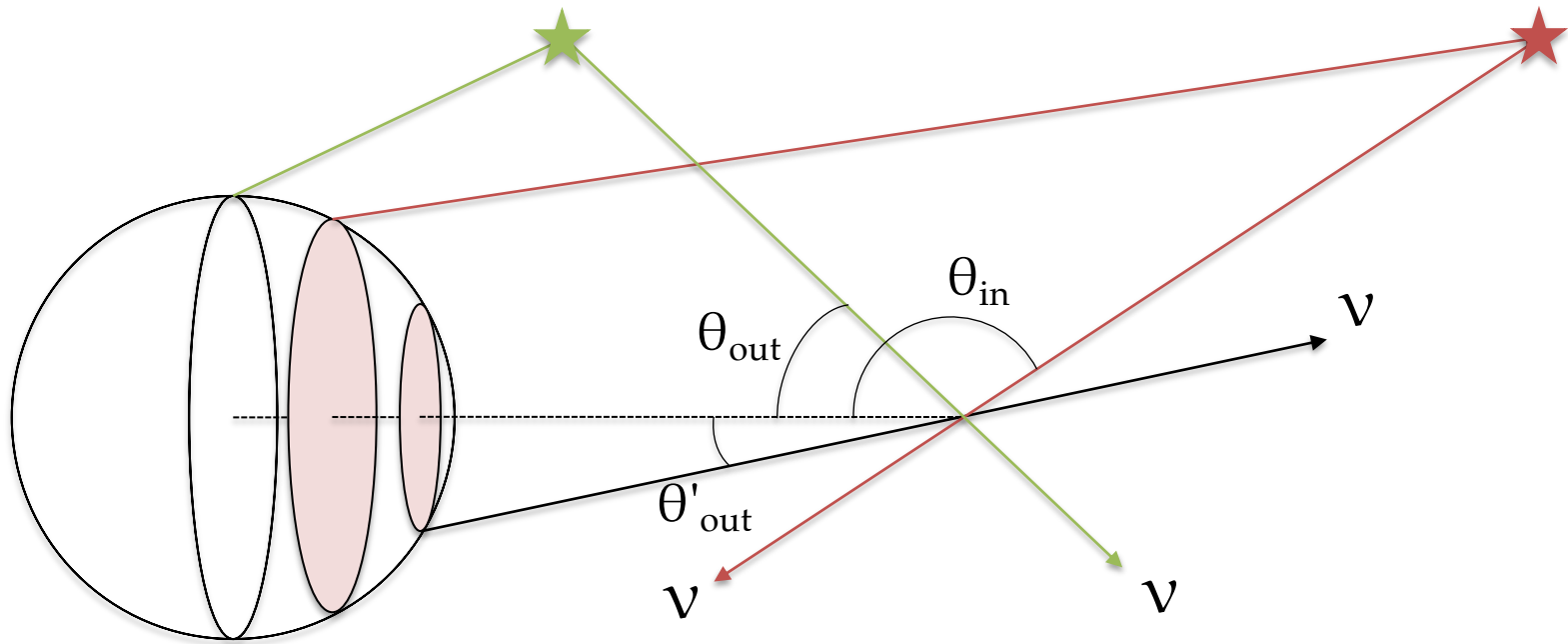
# Inward neutrino flux

We employ iron core-collapse SN model with  $9.6M_{\odot}$ , not O-Ne-Mg SNe.

- Need to expand neutrino halo effects.
- What distribution do inward-going neutrinos have?
- How situation can we ignore them?

We compare inward-going contribution with outward-going one.

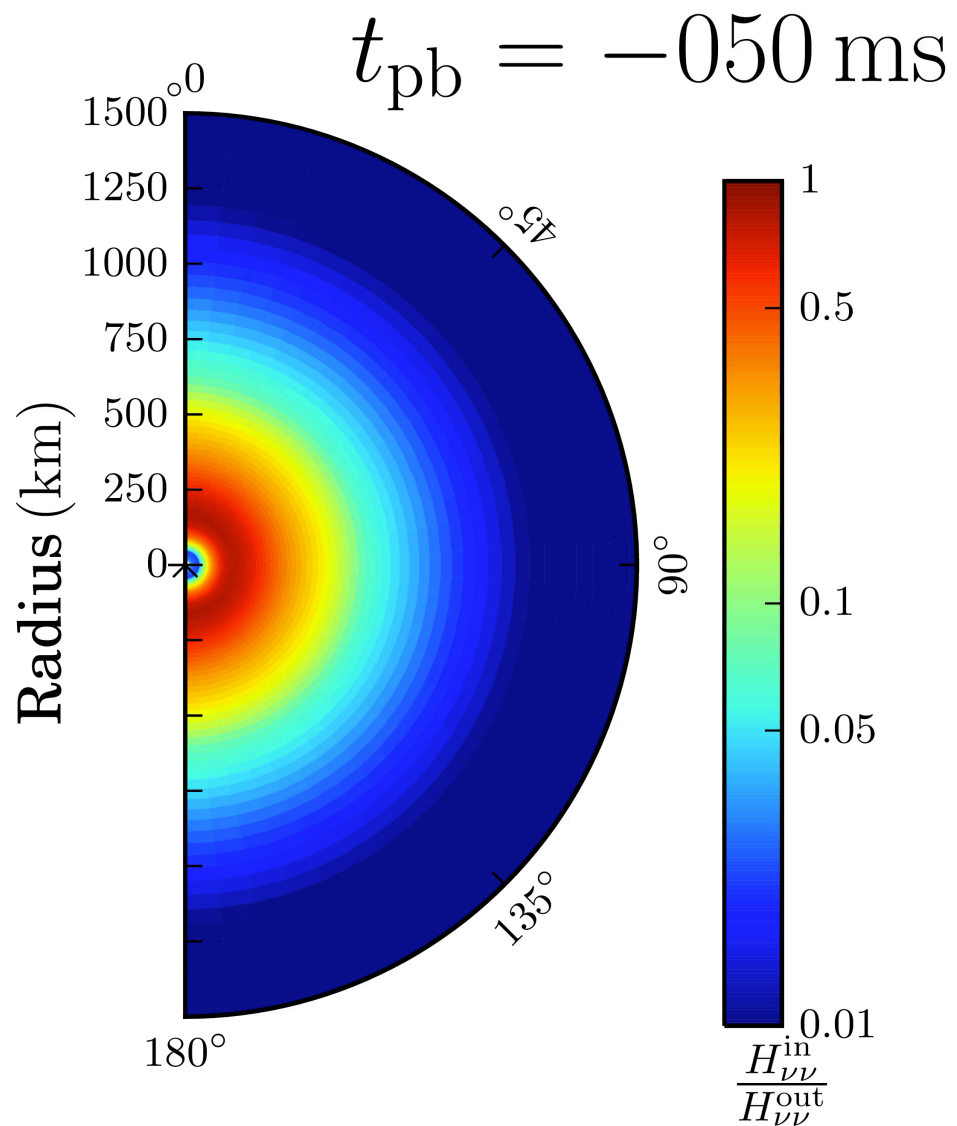
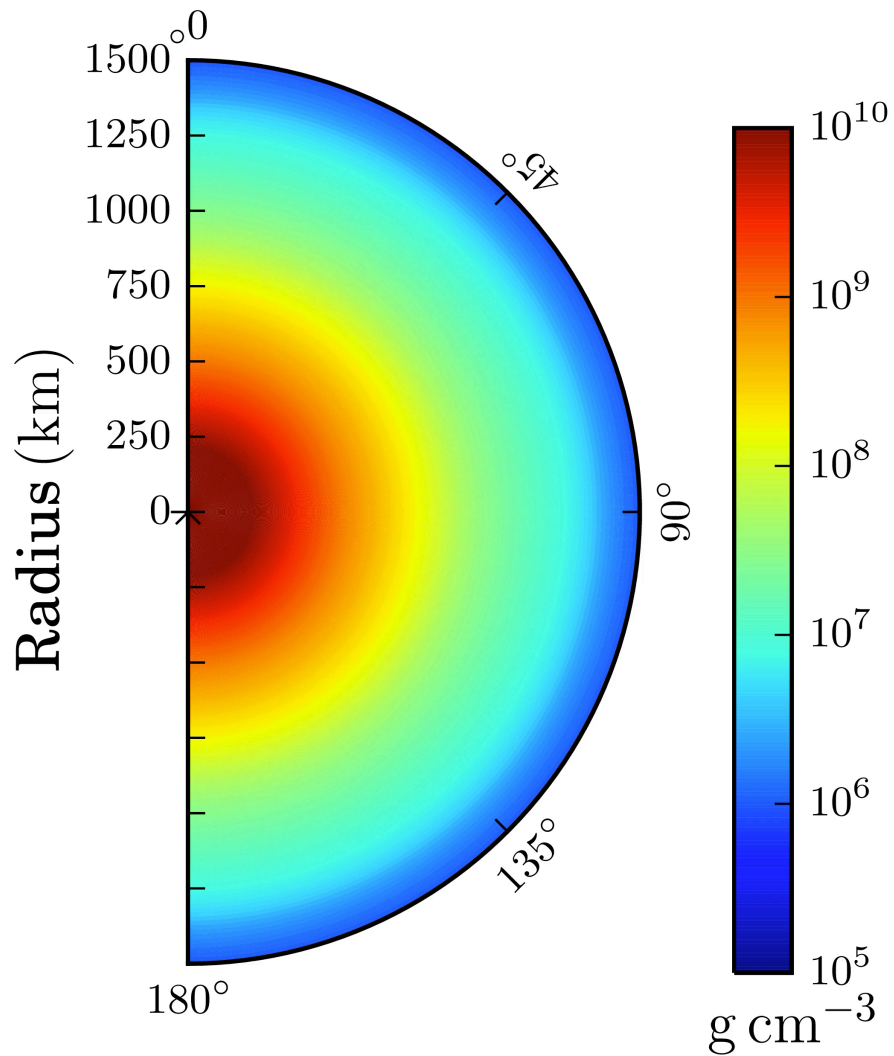
# Inward neutrino flux



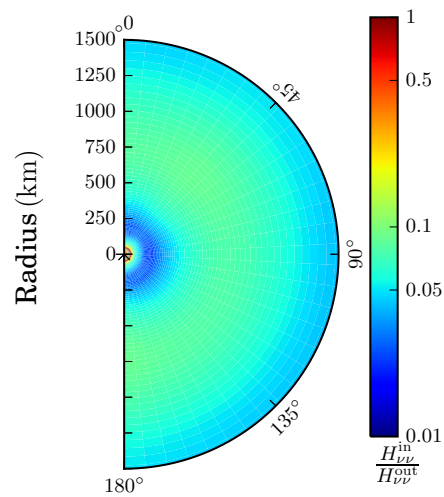
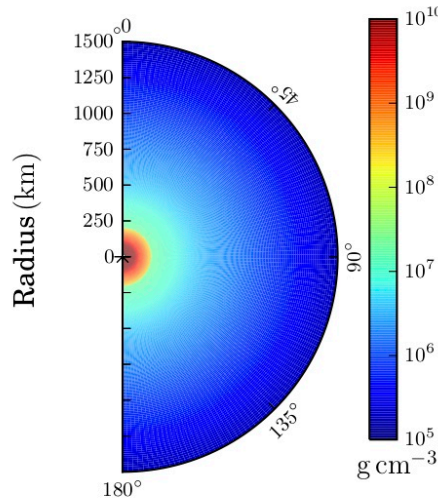
Outward-propagating neutrinos vs. inward-scattered neutrinos

If contribution ratio  $H_{\text{in}}/H_{\text{out}}$  is sufficiently small, we can ignore the inward scattered flux.

$$\langle 1 - \cos \theta_{\text{out}} \rangle \Phi_{\nu_{\text{out}}} \text{ vs. } \langle 1 - \cos \theta_{\text{in}} \rangle \Phi_{\nu_{\text{in}}}$$



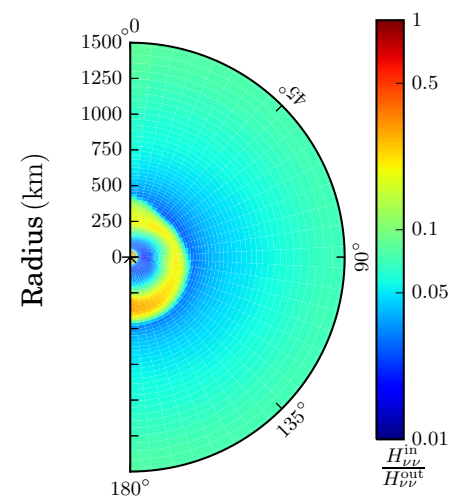
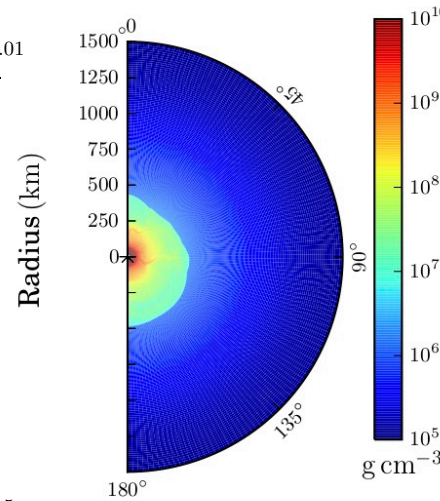
At 86 ms



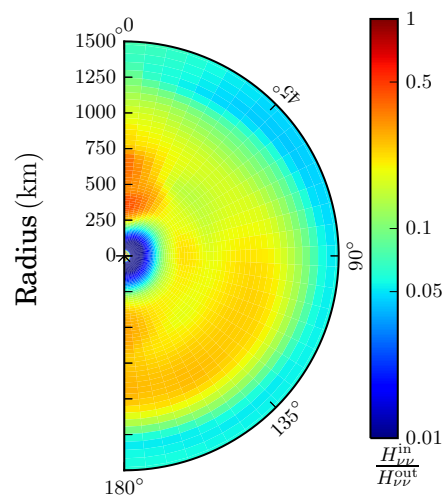
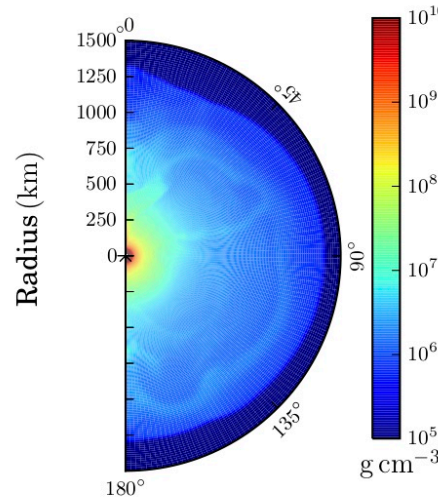
Density profile (left panels)  
Halo distribution (right panels)

Inward scattering is stronger inside the shock wave.

At 136 ms

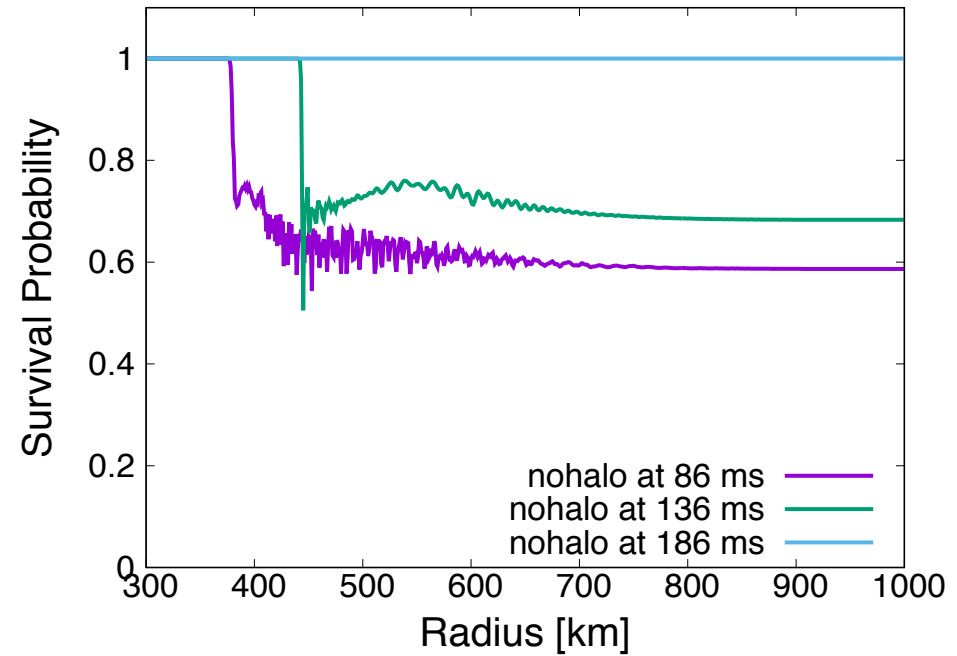
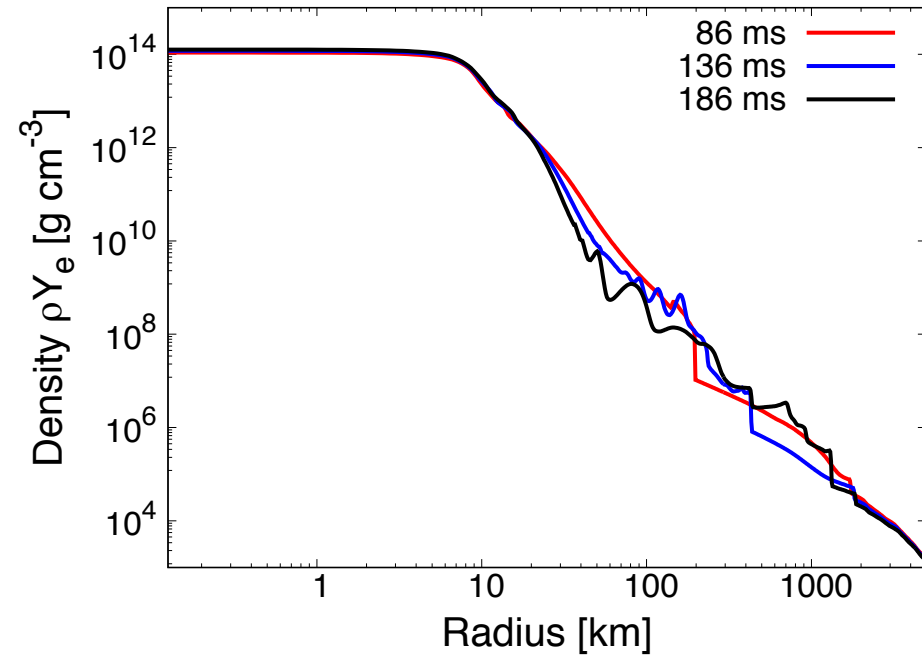


At 186 ms



$$R_{\text{CNO}} = \mathcal{O}(10^2) \text{ km}$$

# CNO w/o halo

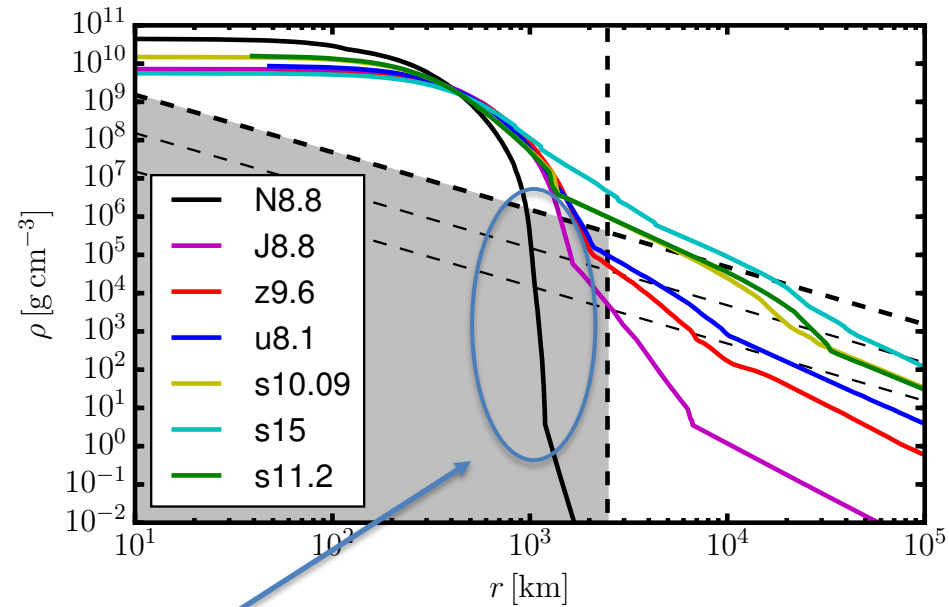
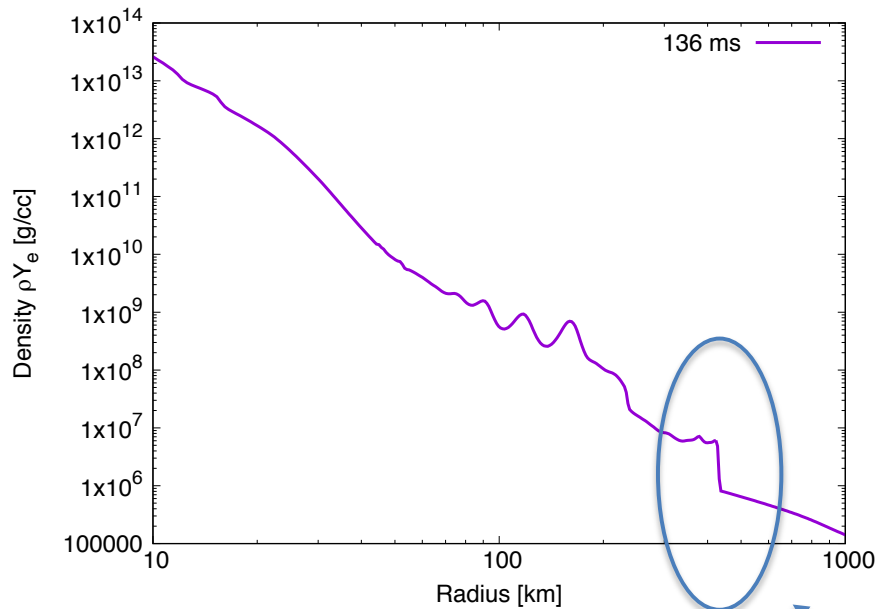


The onset radius of CNO depends strongly on the location of shock wave.

- At 136 ms, CNO starts just after the shock radius.
- At 186 ms, CNO is completely suppressed by high matter density.

Both CNO & halo scattering get free from matter after the shock wave.

# CNO w/o halo

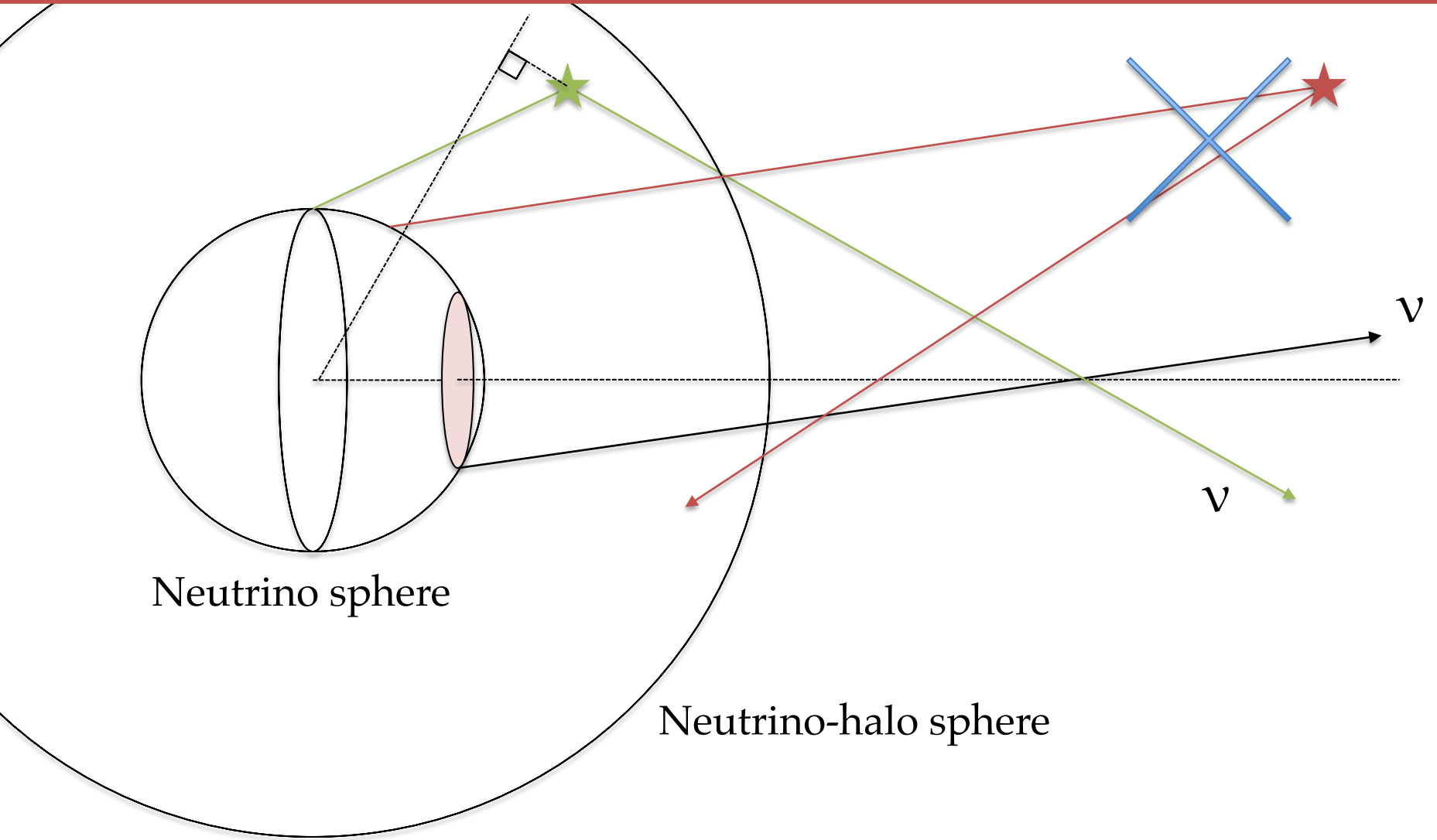


The shock wave in iron-CCSN acts same as  
the steep density gradient in O-Ne-Mg SNe!

We can calculate CNO including neutrino halo safely using BULB-like model.



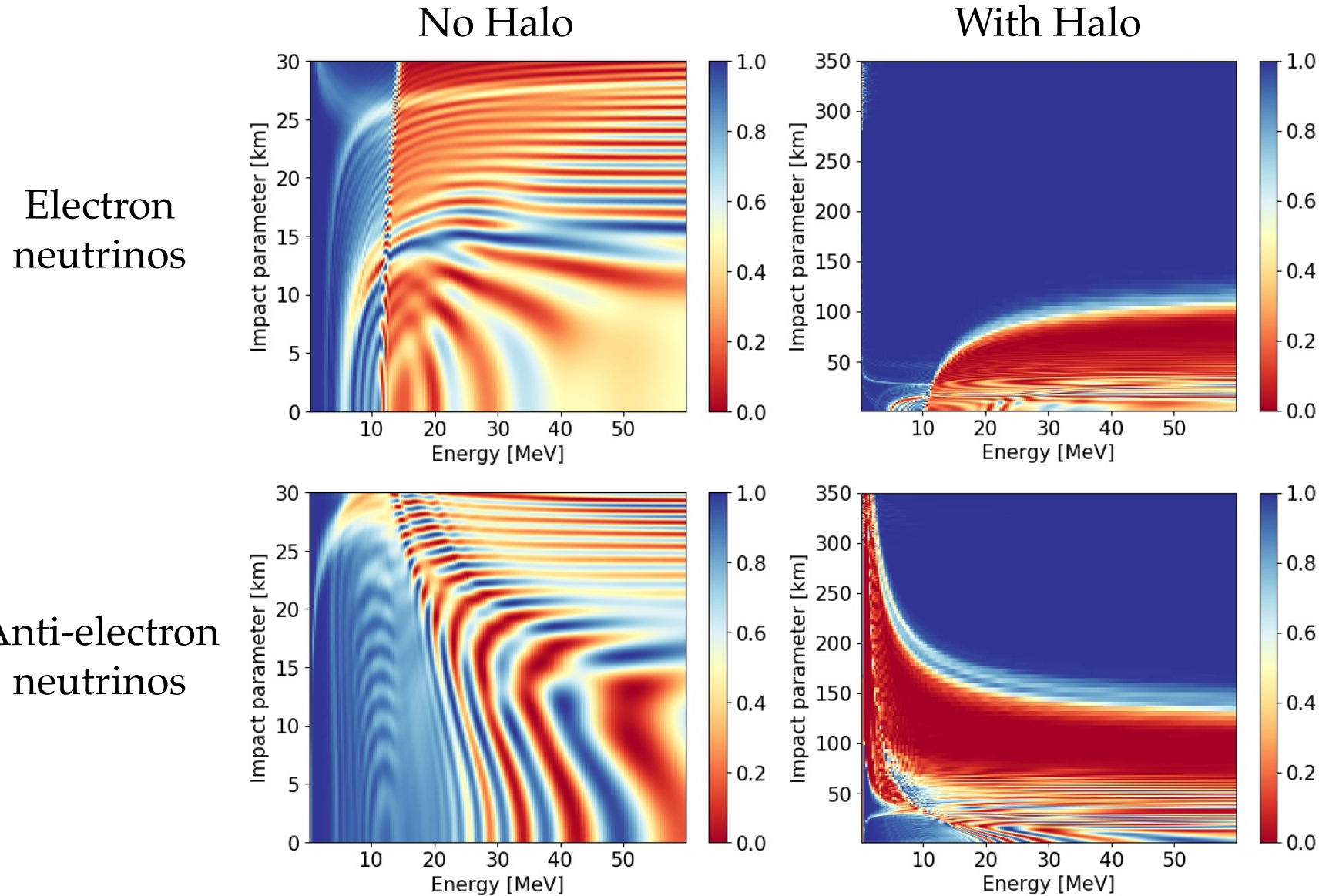
# BULB-like model



Neutrinos are newly considered to be emitted from halo sphere.

# Survival probability after CNO

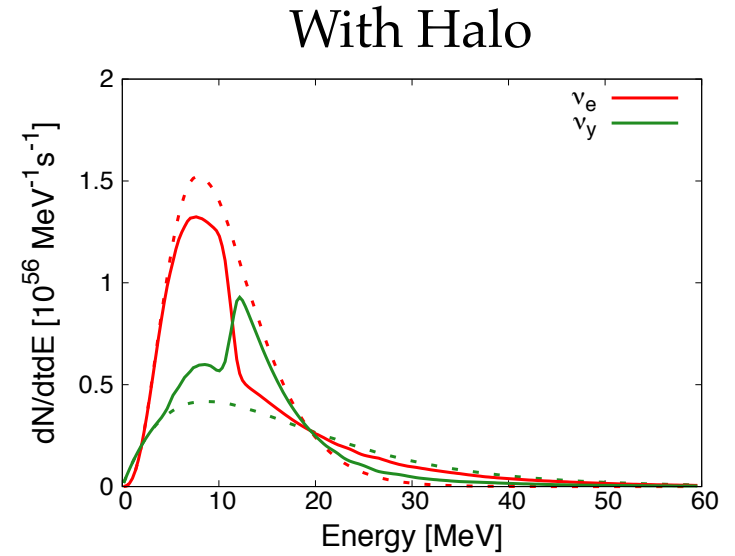
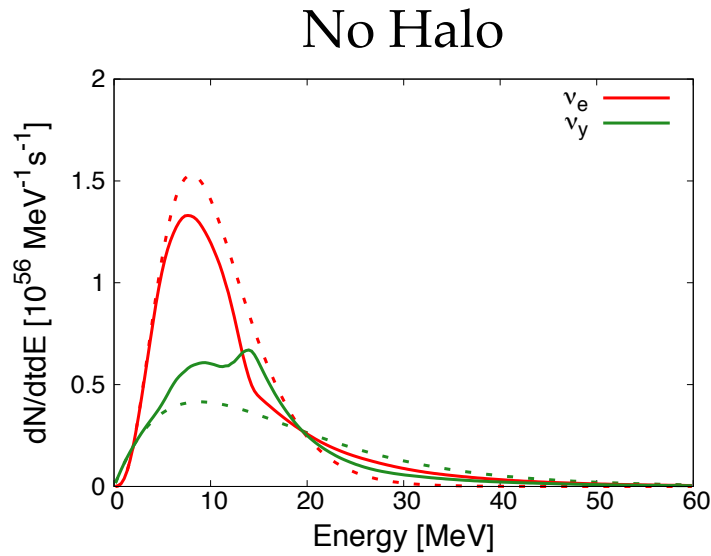
At 136 ms.



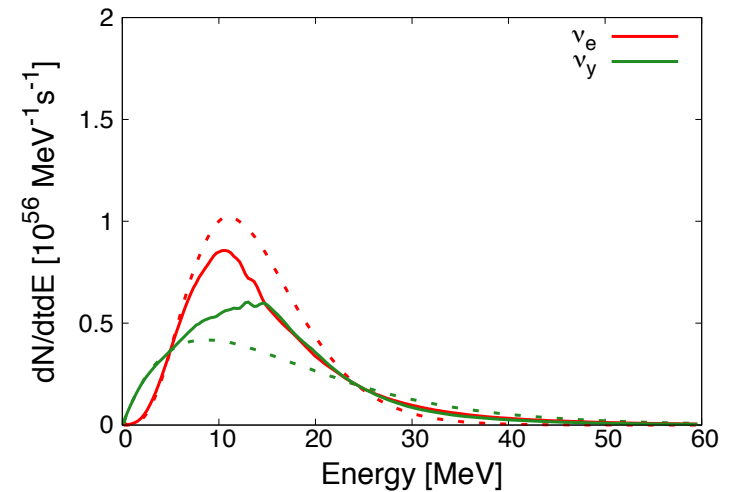
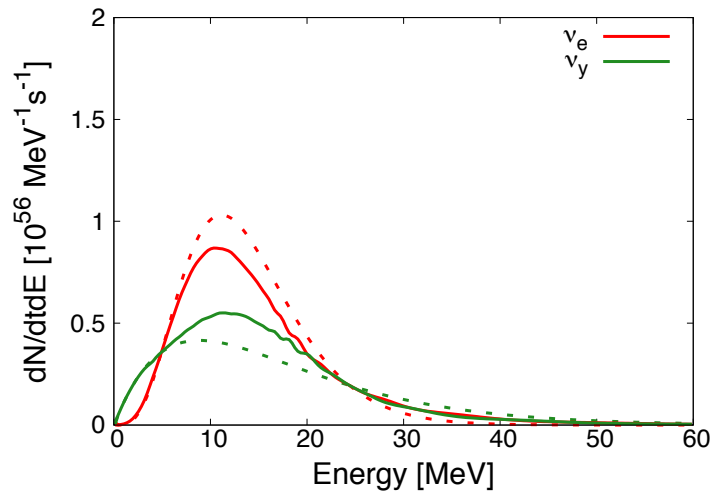
# Spectra after CNO

At 136 ms.

Electron  
neutrinos



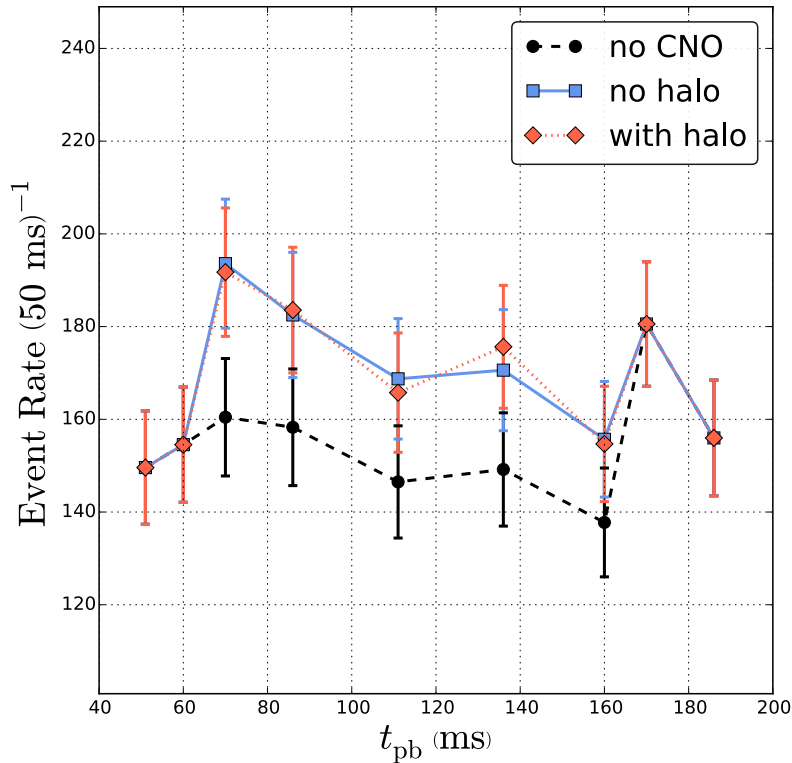
Anti-electron  
neutrinos



# Event rate

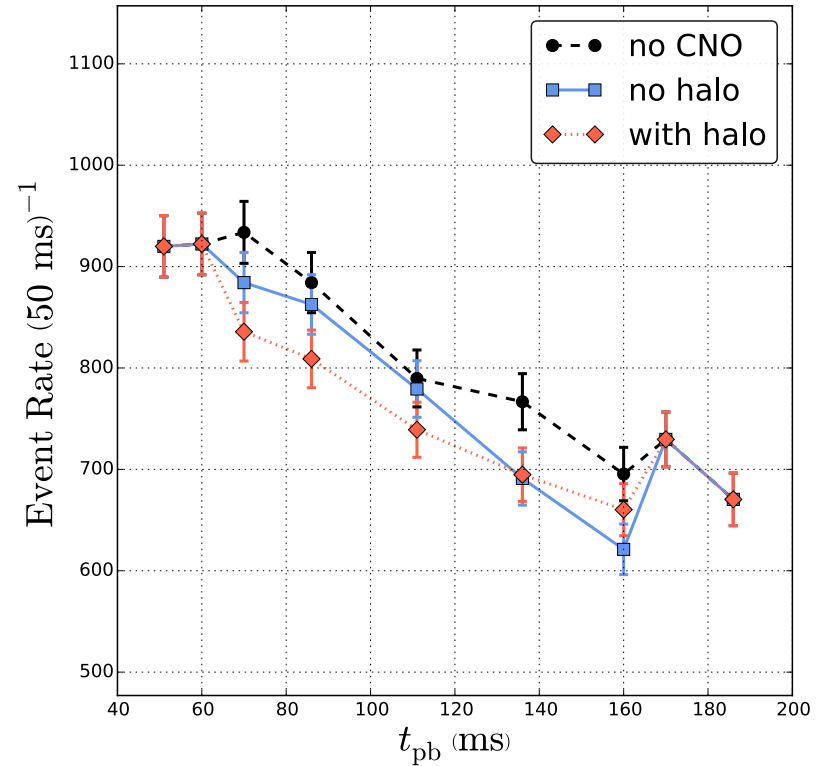
DUNE

Total Event Rate,  $\nu_e - {}^{40}\text{Ar}$ , North polar axis



Super-K

Total Event Rate,  $\bar{\nu}_e - p^+$ , North polar axis



Assume a target at  $d=10\text{kpc}$  in the Galaxy.

Significant differences between no-CNO and inclusion of CNO.

Maximum difference is  $\sim 10\%$  between no-halo and with-halo.

# Conclusion

- BULB model enables us to calculate CNO, but it is not sufficiently to treat the dimensionality problems.
- Recent studies have investigated the beyond BULB model.
  - Some symmetry breakings overcome matter suppression?
- We calculated CNO w/ halo in iron-CCSNe.
  - Inside the shock wave, inward-going neutrinos can not be ignored, compared with outward ones.
  - Outside the shock, the ratio is less than 10%.
  - CNO starts from outside the shock.
- CNO w/ halo has additional flavor conversions over high energy range.